Maximum Caliber Analysis of Ion-Channel Gating

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Introduction

The principle of maximum caliber, MaxCal, is a generalization to nonequilibrium statistical mechanics of the principle of maximum entropy, MaxEnt. E. T. Jaynes introduced the MaxEnt approach to equilibrium statistical mechanics in 1957 [1] and its MaxCal generalization in 1980 [2]. MaxCal has recently been used to derive dynamical laws of transport, analyze single particle two-state dynamics, and study few state models of non-equilibrium processes. [3 - 7] We use MaxCal to analyze hidden Markov models of ion-channel gating and make logical inferences concerning the underlying dynamics. MaxCal is used to determine model parameters; test the adequacy of a model; and predict unmeasured quantities from the trajectory probability distribution.

The Principle of Maximum Caliber

The principle of maximum caliber is to non-equilibrium systems as the principle of maximum entropy is to equilibrium systems.

Equilibrium	
Principle of Maximum Entropy	

Let p_i be the probability that the system is in state i.

Maximize the entropy

$$S = -k_B \sum_i p_i \log(p_i)$$

subject to the constraints

$$\sum_{i} p_{i} = 1$$

$$\langle A_{m} \rangle = \sum_{i} p_{i} A_{mi}$$

Non-Equilibrium Principle of Maximum Caliber

Let p_{Γ} be the probability that the system follows trajectory Γ , where a trajectory is a time sequence of states

Maximize the caliber

$$C = -k_B \sum_{\Gamma} p_{\Gamma} \log(p_{\Gamma})$$

subject to the constraints

$$\sum_{\Gamma} p_{\Gamma} = 1$$

$$\langle A_m(t) \rangle = \sum_{\Gamma} p_{\Gamma} A_{m\Gamma}(t)$$

where the $\langle A_m \rangle' s$ are known expectation values. where the $\langle A_m \rangle' s$ are known expectation values.

The probabilities can be used to determine unknown expectation values

$$\langle B_m \rangle = \sum_i p_i B_{mi}$$

The partition function

$$Z = \sum_{i} \exp(-\sum_{m} \lambda_{m} A_{m})$$

where the λ 's are Lagrange multipliers, can be used to calculate expectation values

$$\langle A_m \rangle = -\frac{1}{Z} \frac{\partial Z}{\partial \lambda_m}$$

The probabilities can be used to determine unknown expectation values

$$\langle B_m(t) \rangle = \sum_{\Gamma} p_i B_{m\Gamma}(t)$$

The dynamic partition function

$$Z_d(T) = \sum_i \exp(-\sum_m \int_0^T \lambda_m(t) A_{m\Gamma}(t) dt)$$

where the λ 's are Lagrange multipliers, can be used to calculate expectation values

$$\langle A_m \rangle = -\frac{1}{Z_d} \frac{\partial Z_d}{\partial \lambda_m}$$

The probabilities calculated using the principle of maximum caliber are the least biased probabilities consistent with the constraints. Observed $\langle A_m \rangle' s$ are used to make unbiased inferences concerning the values of the $\langle B_m \rangle' s$. This point of view is faithful to the original spirit of the principle of maximum entropy put forward by Jaynes in 1957 [1].

Ion Channel Gating

An Ion channel is a membrane protein that can be open, letting a particular species of ion pass through, or closed, blocking the passage of ions. The gating can respond to external voltage, solute concentration, pressure, ligand interactions, etc.



NaK potassium channel, (Bacillus cereus) [8]



5 microseconds of patch-clamp data for an inositol triphosphate receptor [9]

Two-State Model

$$\begin{array}{ccc} & \gamma_{ab} \\ C_a & \leftrightarrows & O_b \\ & \gamma_{ba} \end{array}$$

State a is a closed state, state b is an open state and γ_{ab} is a transition probability.



An example of a two-state trajectory.

The caliber for an N-step time interval is calculated by summing over all possible N-step trajectories. Observables such as the average number of closed state to open state transitions in N time steps, $\langle N_{10} \rangle$, can be measured using single channel gating data. These observables are used to constrain the trajectory probability distribution during the maximization of caliber. Following Ghosh et. al. [3], [4], [5], the partition function when the system is initially in state a can be written as

$$Z_d(N) = (1 \quad 1)G^N \begin{pmatrix} 1\\ 0 \end{pmatrix}$$

where

$$G = \begin{bmatrix} \gamma_{aa} & \gamma_{ab} \\ \gamma_{ba} & \gamma_{bb} \end{bmatrix}$$

and N is the number of time steps in the trajectories.

Two-State Results



The expected fractional number of state a to state b transitions verses the number of time steps with $\gamma_{ab} = 0.5$, $\gamma_{ba} = 0.6$, and an initial state a. The solid line is the MaxCal prediction and the points are Monte Carlo simulation results.

Hidden States

$$\begin{array}{ccc} \gamma_{ab} & \gamma_{bc} \\ C_a \leftrightarrows C_b \leftrightarrows O_c \\ \gamma_{ba} & \gamma_{cb} \end{array}$$

In this simple model there are two closed states, one open state, and transitions between states a and c are not allowed. Since a closed channel can either be in state a or state b we call these hidden states. We solved the problem of applying MaxCal to a model with hidden states by considering multiple-time-step observables such as $\langle N_{000} \rangle$, the expected number of times the system will make a transition from a closed state to another closed state and then to another closed state.



A three-state trajectory showing a time sequence of transitions between states a, b and c and the corresponding sequence of open and closed transitions.

The partition function for this system can be written as

$$Z_{d}(N) = (1 \quad 1 \quad 1)G^{N}\begin{pmatrix} p_{a} \\ p_{b} \\ 0 \end{pmatrix} \qquad G = \begin{bmatrix} \gamma_{aa} & \gamma_{ab} & 0 \\ \gamma_{ba} & \gamma_{bb} & \gamma_{bc} \\ 0 & \gamma_{cb} & \gamma_{cc} \end{bmatrix}$$

where N is the number of time steps in a trajectory and p_a and p_b are the equilibrium occupation probabilities for states a and b when the system is initially in a closed state.

Hidden-State Results

The results of our analysis with

$$G = \begin{bmatrix} 0.7 & 0.2 & 0 \\ 0.3 & 0.5 & 0.3 \\ 0 & 0.3 & 0.7 \end{bmatrix}$$

are given in the table below. We first tried to analyze the simulated three-state data with a twostate model. The values of $\langle N_{10} \rangle$ and $\langle N_{01} \rangle$ were used to fix the parameters of the model then the observables $\langle N_{000} \rangle$ and $\langle N_{0001} \rangle$ were calculated. The disagreement between the actual values for these observables and the two-state model prediction is a flag that tells us that two states are not enough. When the correct number of states were used, MaxCal gave us the correct model parameters.

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Observable	Simulation	Two-State Results	Three-State Model Results
$\langle N_{10} \rangle$	0.180 ± 0.001	0.180 ± 0.001	0.180
$\langle N_{01} \rangle$	0.300 ± 0.001	0.300 ± 0.001	0.300
(N ₀₀₀)	0.694 ± 0.002	0.672 ± 0.001	0.694
(N ₀₀₀₁)	0.166 ± 0.001	0.121 ± 0.001	0.165

Observables calculated from the MaxCal analysis of simulated three-state data. Note the disagreements (bold).

Conclusions

For Markov models of Ion-channel gating, including hidden Markov models, the principle of maximum caliber can be used to:

- 1. determine model parameters,
- 2. test the adequacy of the model, (i.e. Do we need to add states?), and
- 3. predict unmeasured quantities from the trajectory probability distribution.

References

- 1. E.T. Jaynes, Phys. Rev. 106, 620 (1957).
- 2. E.T. Jaynes, Ann. Rev. Phys. Chem. 31, 579 (1980).

3. K. Ghosh, K. Dill, M. Inamdar, E. Seitaridou, and R. Phillips, Am. J. Phys. 74, 123 (2006).

4. Stock, G., K. Ghosh, and K. Dill, J. Chem. Phys. 128, 194102 (2008).

5. D. Wu, K. Ghosh, M. Inamdar, H.J. Lee, S. Fraser, K. Dill, and R. Phillips, Phys. Rev. Lett. 103, 050603 (2009).

- 6. M. Otten and G. Stock, Phys. Rev. E 82, 031905 (2010).
- 7. S. Presse, K. Ghosh, R. Phillips, and K. Dill, Phys. Rev. E 82, 031905 (2010).
- 8. M. Lomize, A. Lomize, I. Pogozheva, H. Mosberg, Bioinformatics 22, 623 (2006).
- 9. E. Gin, M. Falcke, L. E. Wagner, D.I. Yule, J. Sneyd, J. Theo. Biol. 257, 460 (2009).